Chapter 29 Electric Currents and Magnetic Fields

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Biot & Savart's Experiments

History

1820 AD – Hans Christian Ørsted (Danish) Electric current generates magnetic fields.

1820 AD (18 December) – Jean-Baptiste Biot (French) & Félix Savart (French) Experimental determination of current generated magnetic fields. Treat currents as magnets with magnetic fields.

$$d\vec{B} = \frac{\mu_0}{4\pi} \frac{Id\vec{l} \times \hat{r}}{r^2} \qquad \frac{\mu_0}{4\pi} = 10^{-7} T m / A$$

1820 AD (4 December) – André-Marie Ampère (French) interaction between currents

1882 AD – William Thomson, Baron Kelvin (Scottish) current balance design



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The Biot-Savart Law

Magnetic Fields

Use the Biot-Savart law to calculate the magnetic field at the center of a circular loop with radius *R* and current *I*. Please calculate the magnetic field on the central axis and at a distant *d* away from the center.

$$\vec{B}_{o} = \frac{\mu_{0}I}{4\pi} \int \frac{d\vec{l} \times (-\hat{r})}{r^{2}} = \frac{\mu_{0}I}{4\pi} \int_{0}^{2\pi} \frac{Rd\theta}{R^{2}} \hat{k}$$

$$\vec{B}_{o} = \frac{\mu_{0}I}{4\pi} \frac{2\pi}{R} \hat{k} = \frac{\mu_{0}I}{2R} \hat{k}$$

$$\vec{B}_{d} = \frac{\mu_{0}I}{4\pi} \int \frac{d\vec{l} \times (-\hat{r})}{r^{2}} = \frac{\mu_{0}I}{4\pi} \int_{0}^{2\pi} \frac{Rd\theta}{R^{2} + d^{2}} \frac{R}{\sqrt{R^{2} + d^{2}}} \hat{k}$$

The Biot-Savart Law

Magnetic Fields

Use the Biot-Savart law to calculate the magnetic field at the center of
a circular loop with radius *R* and current *I*. Please calculate the
magnetic field on the central axis and at a distant *d* away from the
center.
$$\vec{B}_{d} = \frac{\mu_{0}I}{4\pi} \frac{R^{2}2\pi}{(R^{2} + d^{2})^{\frac{3}{2}}} \hat{k} = \frac{\mu_{0}IR^{2}}{2(R^{2} + d^{2})^{\frac{3}{2}}} \hat{k}$$
$$\vec{F}_{d} = \frac{R}{4\pi\epsilon_{0}} \frac{2p}{r^{3}}$$
$$B_{d} \xrightarrow{\mu_{0}IR^{2}}{2d^{3}} = \frac{2\mu_{0}\pi IR^{2}}{4\pi d^{3}} = \frac{\mu_{0}}{4\pi} \frac{2m}{d^{3}} \iff E = \frac{1}{4\pi\epsilon_{0}} \frac{2p}{r^{3}}$$
$$d\vec{B} = \frac{\mu_{0}}{4\pi} \frac{Id\vec{l} \times \hat{r}}{r^{2}}$$

The Biot-Savart Law

Magnetic Fields

Use the Biot-Savart law to calculate the magnetic field at point P a distance *a* away from the segment of line current *I*.

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$$\vec{B}_{P} = \frac{\mu_{0}I}{4\pi} \int \frac{d\vec{l} \times \hat{r}}{r^{2}} = \frac{\mu_{0}I}{4\pi} \int_{\theta_{1}}^{\theta_{2}} \frac{\sin\left(\frac{\pi}{2} - \theta\right)}{(a/\cos(\theta))^{2}} d(a\tan(\theta)) \hat{k} \qquad x = a\tan\theta$$

$$\vec{B}_{P} = \frac{\mu_{0}I}{4\pi} \int_{\theta_{1}}^{\theta_{2}} \frac{\cos(\theta) a}{a^{2}} d\theta \hat{k}$$

$$\vec{B}_{P} = \frac{\mu_{0}I}{4\pi a} (\sin(\theta_{2}) - \sin(\theta_{1})) \hat{k}$$
An infinitely long line current, $\theta_{2} = \frac{\pi}{2}, \theta_{1} = -\frac{\pi}{2}$

$$\vec{B}_{P} = \frac{\mu_{0}I}{2\pi a} \hat{k}$$

Ampere's Law & Applications for The Calculation of Magnetic Fields

Magnetic fields make a circular rotation around the current lead. The magnitude of the magnetic field is of the same value of $\mu_0 I/2\pi r$. Line integral of the field gives

$$\oint \vec{B} \cdot d\vec{l} = \int_0^{2\pi} \frac{\mu_0 I}{2\pi r} r d\theta = \mu_0 I$$

Ampere's Law:

Ampere's Law

$$\oint \vec{B} \cdot d\vec{l} = \mu_0 I$$
$$\oint \vec{B} \cdot d\vec{l} = \mu_0 I_{enclosed}$$



There Are No Magnetic Monopoles.

Gauss's Law for Magnetic Fields Magnetic moment (magnetization) possesses the same concept as magnetic dipoles.

In analogy to electric dipole, the Gauss's law gives us zero magnetic charges in a box due to the same number of m + and m - .



$$\oint \vec{E} \cdot d\vec{A} = \frac{1}{\varepsilon_0} \left(+Nq + (-Nq) \right) = 0$$

$$\oint \vec{B} \cdot d\vec{A} = \mu_0 \left(+Nm + (-Nm) \right) = 0$$

Ref: https://phys.org/news/2016-08-mysterious-magnetic-monopole.html

Magnetic Force

Use Ampere's law to find the magnetic field around an infinite longe wire carrying current *I*. Use Lorentz law to find the force per unit length between two parallel current-carrying and straight wires.

$$\oint \vec{B} \cdot d\vec{l} = \mu_0 I \rightarrow 2\pi r B = \mu_0 I$$

$$B = \frac{\mu_0 I}{2\pi r}$$

$$F = ILB = \frac{\mu_0 I^2}{2\pi r} L \quad \Longrightarrow \quad \frac{F}{L} = \frac{\mu_0 I^2}{2\pi r}$$
If $r = 1$ m, $I = 1$ A,
$$\frac{F}{L} = 2(10^{-7}) \frac{1^2}{L} = 2 \times 10^{-7} N/m$$

Magnetic Force Between Two Parallel, Current-Carrying Wires



The Magnetic Field of a Solenoid

Solenoid

id

$$\vec{B}_d = \frac{\mu_0 I R^2}{2(R^2 + d^2)^{\frac{3}{2}}} \hat{k}$$
 \hat{x}
 \hat{y}

Let the current be I and the number of coils per unit length n for the solenoid. With a small displacement dx, the coil number is ndx and the current is *Indx*. $d\vec{B} = \hat{\imath} - \frac{\mu_0 R^2}{3} Indx$ $2(R^2 + x^2)^{\overline{2}}$ $B = \frac{\mu_0 n I R^2}{2} \int_{x_1}^{x_2} \frac{dx}{(R^2 + x^2)^{\frac{3}{2}}} \quad \text{Let } x = R \tan \theta$ $B = \frac{\mu_0 n I R^2}{2} \int_0^{\theta_2} \cos \theta \, d\theta = \frac{1}{2} \mu_0 n I (\sin \theta_2 - \sin \theta_1)$

For an infinite long solenoid, $\theta_1 = -\pi/2$ and $\theta_2 = \pi/2$, we get $B = \mu_0 nI$

The Magnetic Field of a Solenoid

Solenoid

Using the Ampere's law to find the magnetic field of a very long solenoid. Consider *n* closely wound turns per unit length on a cylinder of radius R and carrying a steady current I. Turn to see the crosssection. $\bigotimes \bigotimes \bigotimes$ I out of page *I* in For the wires of current on the top, we draw an Ampere's loop and find the magnetic field by $\otimes \otimes_{R} \otimes \otimes \otimes \otimes$ $\int \Delta w \ll l$ $Bl + Bl = \mu_0 n l l$ $B = \frac{1}{2}\mu_0 nI$

Add up the magnetic field due to the wires of current on the bottom:

$$B = \mu_0 n I$$

The Magnetic Field of a Solenoid

Magnetism in Atoms

Estimate the orbital magnetic moment of the electron in an hydrogen atom. $\vec{L} = \vec{r} \times (m\vec{v}) \rightarrow L = mvr$, the quantization rule gives myou $L = n\hbar$ Estimate the current due to the charge -e. $I = -\frac{e}{T} = -\frac{e}{2\pi r} = -\frac{ev}{2\pi r}$ Estimate the magnetic moment. $\mu = IA \rightarrow \mu = -\frac{ev}{2\pi r}\pi r^2 = -\frac{evr}{2} = -\frac{e}{2m}mvr = -\frac{e}{2m}L$ Consider the quantization effect, you have $\mu = -n \frac{e\hbar}{2m}$. Here we define the Bohr magneton $\mu_B = \frac{e\hbar}{2m}$ as a unit for atomic magnetic moments. $\mu = -n\mu_{\rm B}$

Similar to the quantization of orbital angular momentum as:

$$\mu_L = -\frac{e}{2m}L = -n\frac{e\hbar}{2m} = -\frac{\mu_B}{\hbar}n\hbar = -g_L\frac{\mu_B}{\hbar}I$$

Considering the relativistic effect, the Dirac equation results in another quantization, spin quantization, giving another similar equation:

$$\mu_{S} = -g_{S} \frac{\mu_{B}}{\hbar} S$$

The differences are listed below.
$$g_{L} = 1, g_{S} = 2, L = n\hbar, S = \frac{n}{2}\hbar$$

	Fe (BCC)	Co (HCP)	Ni (FCC)
Bohr Magnettons per Atom	2.22	1.72	0.60

Magnetism in Atoms

Magnetism in Matter



The Magnetic Torque

Examples

A magnet of magnetic moment m is placed at the origin on the xzplane with a direction making an angle of $\theta = 30^{\circ}$ to the x-axis. A coil of 12 loops is fixed at the origin with its central axis along the x-axis. Please calculate the magnetic torque on the magnet. y



Calculation of The Magnetic Field

Find the magnetic field at the center of a square loop of edge length L, which carries a current I.

Use the previously derived equation with $\theta_1 = -\pi/4$ and $\theta_2 = \pi/4$ for one edge of the square loop. $B = \frac{\mu_0 I}{4\pi r} (\sin(\theta_2) - \sin(\theta_1))$ $B_{one \ edge} = \frac{\mu_0 I}{4\pi (L/2)} (\sin(\pi/4) - \sin(-\pi/4)) = \frac{\mu_0 I}{2\pi L} \sqrt{2}$ $B_{total} = 4 \frac{\mu_0 I}{2\pi L} \sqrt{2} = \frac{2\mu_0 I}{\pi L} \sqrt{2}$



The Magnetic Torque

Examples

Two infinite wires carrying current *I* are placed on the xz plane with their central axes parallel to the *z*-axis. They are symmetrically arranged on both sides of the *z*-axis with a separation of 2d. Please calculate the magnetic field at the position of (0, y, 0).

$$B_{wire} = \frac{\mu_0 I}{2\pi r} \qquad r = \sqrt{d^2 + y^2}$$

$$\vec{B}_{1} = \frac{\mu_{0}I}{2\pi\sqrt{d^{2} + y^{2}}} \left(-\frac{\sqrt{2}}{2}\hat{\imath} - \frac{\sqrt{2}}{2}\hat{\jmath}\right)$$

$$\vec{B}_{total} = \frac{\mu_0 I}{2\pi\sqrt{d^2 + y^2}} \left(-\sqrt{2}\hat{\imath}\right)$$



Two straight rods of *L* in length are separated at a distance *d* apart in a current balance. The carry currents are *I* each in opposite direction. What mass must be placed on the upper rod?

The magnetic field due to an infinite wire is

 $B = \frac{\mu_0 I}{2\pi d}$

The repulsive Lorentz force is

$$F = IlB = IL \frac{\mu_0 I}{2\pi d}$$
$$mg = \frac{\mu_0 I^2 L}{2\pi d} \rightarrow m = \frac{\mu_0 I^2 L}{2\pi g d}$$



Current & Current Density:

Ampere's Law & Applications for The Calculation of Magnetic Fields

$$I = JA \rightarrow J = \frac{I}{A} = \frac{I}{\pi R^2}$$

$$r > R: \int \vec{B} \cdot d\vec{l} = B2\pi r = \mu_0 I \rightarrow B = \frac{\mu_0 I}{2\pi r}$$

$$r < R: \int \vec{B} \cdot d\vec{l} = \mu_0 I \rightarrow 2\pi r B = \mu_0 (J\pi r^2)$$

$$B = \frac{\mu_0 J r}{2} = \frac{\mu_0 I r}{2\pi R^2}$$

Ampere's Law & Applications for The Calculation of Magnetic Fields

A wire carrying a current *I* is bent to form a toroid with *N* loops. Please find the magnetic field inside and outside the toroid.

NTT

Outside the toroid:

$$r > b: 2\pi rB = \mu_0 \times 0 \rightarrow B = 0$$

$$r < a: 2\pi r B = \mu_0 \times 0 \rightarrow B = 0$$

Inside the toroid:

$$b > r > a: 2\pi rB = \mu_0 \times NI \rightarrow B = \frac{\mu_0 NI}{2\pi r}$$



Ampere's Law & Applications for The Calculation of Magnetic Fields

Examples

A very long, thin strip of metal of width *w* carries a current *I* along its length. The current is distributed uniformly across the width of the strip. Find the magnetic field at point P in the diagram. Point P is in the plane of the strip at distance *b* away from its edge.

Use Ampere's law

$$2\pi y(dB) = \mu_0 \frac{I}{w} dy$$

$$B = \int_{b}^{b+w} \frac{\mu_0 I}{2\pi y w} dy$$

 $B = \frac{\mu_0 I}{2\pi w} \ln\left(\frac{b+w}{b}\right)$



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